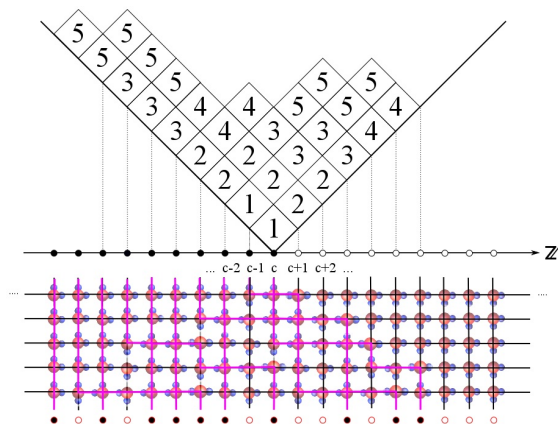


The asymmetric six-vertex model & cylindric Hecke characters

Based on COMM. MATH. PHYS. 381 (2021) 591-640



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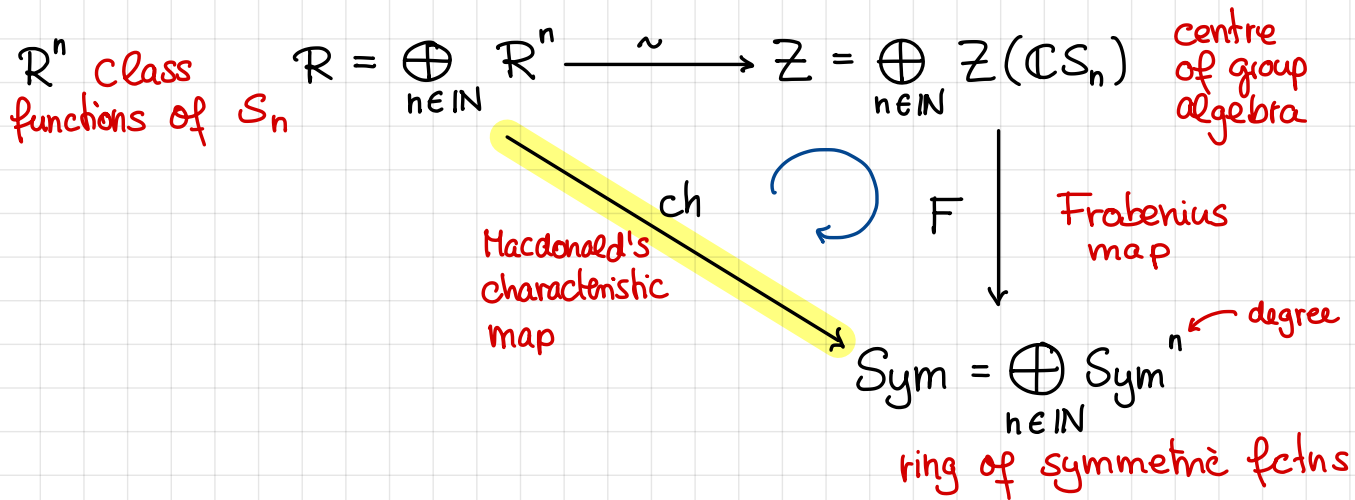
University
of Glasgow

Exactly solvable lattice models, Stanford, 8 Mar 2022

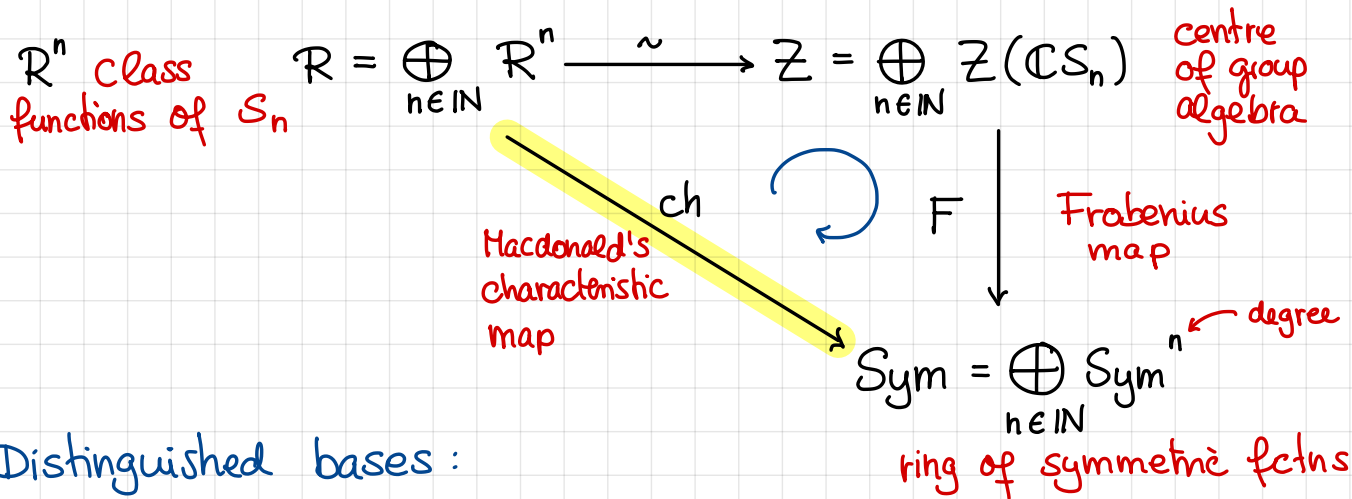
Outline

- ① Madonald's characteristic map & the Boson-Fermion Correspondence
- ② Hecke characters & the asymmetric six-vertex model
- ③ Vertex operators
- ④ Periodic boundary conditions: cylindric Hecke characters
& Gromov-Witten invariants

§1. Macdonald's characteristic map & the BF correspondence



§1. Macdonald's characteristic map & the BF correspondence



Distinguished bases:

- ① \mathbb{R}^n : irreducible characters χ^λ (Specht module), $\lambda + n$
- ② $\mathbb{Z}(\mathbb{C}S_n)$: sums over conjugacy classes $c_\lambda = \sum_{\text{cycle type } \lambda} w$
- ③ Sym^n : (a) power sums $\{p_\lambda : \lambda + n\}$
 (b) Schur functions $\{s_\lambda : \lambda + n\}$

§1. Macdonald's characteristic map & the BF correspondence

$$\begin{array}{ccc}
 \mathbb{R} = \bigoplus_{n \in \mathbb{N}} \mathbb{R}^n & \xrightarrow{\sim} & \mathbb{Z} = \bigoplus_{n \in \mathbb{N}} \mathbb{Z}(\mathbb{C}S_n) \\
 \searrow \text{ch}(\chi^\lambda) = s_\lambda & & \downarrow F(c_\mu) = \frac{p_\mu}{z_\mu} \\
 & & \text{Sym} = \bigoplus_{n \in \mathbb{N}} \text{Sym}^n
 \end{array}$$

power sums

Schur function

THM (Macdonald) The above diagram of vector space isomorphisms commutes.

COR (Frobenius formula)
$$s_\lambda = \sum_{\mu} \frac{\chi_\lambda(\mu)}{z_\mu} p_\mu$$

inverse
 Schur function \nearrow s_λ \leftarrow $\chi_\lambda(\mu)$ \leftarrow p_μ \leftarrow power sum
 \nearrow cycle type \leftarrow size of conjugacy class

§ 1. Macdonald's characteristic map & the BF correspondence

$$\mathbb{R} = \bigoplus_{n \in \mathbb{N}} \mathbb{R}^n \xrightarrow{\sim} \mathbb{Z} = \bigoplus_{n \in \mathbb{N}} \mathbb{Z}(\mathbb{C}S_n)$$

$$\begin{array}{ccc} & & \downarrow \text{Frob} \\ & \searrow \text{ch} & \\ & \text{Macdonald's} & \text{Frobenius} \\ & \text{characteristic} & \text{map} \\ & \text{map} & \\ & & \downarrow \\ & & \text{Sym} = \bigoplus_{n \in \mathbb{N}} \text{Sym}^n \end{array}$$

Hopf algebra structure :

$$\text{Ind}_{S_m \times S_n}^{S_{m+n}} : \mathbb{R}^m \otimes \mathbb{R}^n \rightarrow \mathbb{R}^{m+n}$$

$$\text{Res}_{S_m \times S_n}^{S_N} : \mathbb{R}^N \rightarrow \bigoplus_{m+n=N} \mathbb{R}^m \otimes \mathbb{R}^n$$

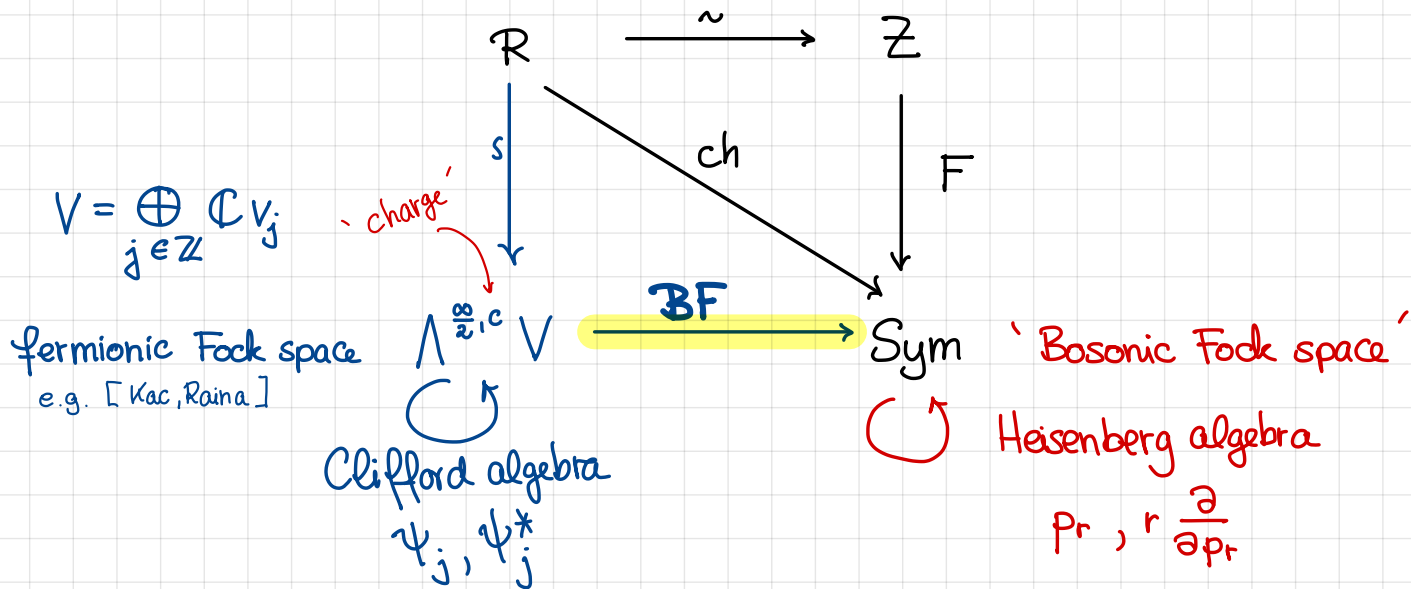
$c_{\lambda\mu}^\nu =$ LR coefficients

$$s_\lambda \cdot s_\mu = \sum_\nu c_{\lambda\mu}^\nu s_\nu$$

$$\Delta s_\lambda = \sum_{\mu \subset \lambda} s_{\lambda/\mu} \otimes s_\mu$$

$$s_{\lambda/\mu} = \sum_\nu c_{\mu\nu}^\lambda s_\nu$$

§1. Macdonald's characteristic map & the BF correspondence



THM The BF correspondence is an Heisenberg algebra module isomorphism

$$P_r = \sum_{j \in \mathbb{Z}} \psi_j \psi_{j+r}^* \mapsto p_r, \quad P_r^* = \sum_{j \in \mathbb{Z}} \psi_{j+r} \psi_j^* \mapsto r \frac{\partial}{\partial p_r}$$

The fermionic Fock space

Identify elements in $\bigwedge^{\infty} V$ as \bigvee (half-) infinite wedge products the span of

$$V = V_{i_1} \wedge V_{i_2} \wedge V_{i_3} \wedge \dots \quad \text{with } i_k = c - k \text{ for } k \ll -1$$

Define a Maya diagram $\sigma_{\lambda, c}: \mathbb{Z} \rightarrow \mathbb{Z}^2$ via $\sigma_{\lambda, c}(j) = \begin{cases} 1, & j = c + 1 + \lambda_i - i \\ 0, & \text{else} \end{cases}$

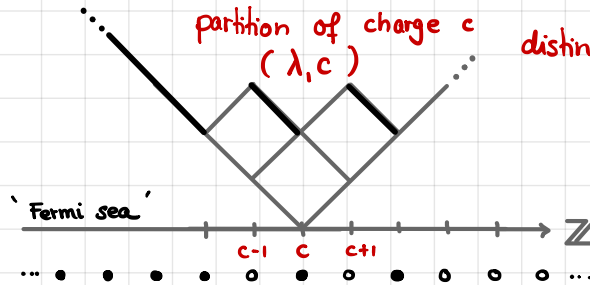
distinguished basis:

Maya diagrams

$$\sigma: \mathbb{Z} \rightarrow \{0, 1\}$$

$$\psi_i \psi_j^* + \psi_j^* \psi_i = \delta_{ij}$$

Clifford algebra



distinguished basis: power sums

$$p_r = \sum_{i \geq 1} x_i^r$$

Heisenberg algebra

$$\sigma_{\lambda, c} \mapsto \langle \sigma_{\lambda, c}, e^{H[P]} \sigma_{\emptyset, c} \rangle = s_{\lambda}[X], \quad H[P] = \sum_{r > 0} \frac{p_r}{r} P_r, \quad P_r = \sum_{i \in \mathbb{Z}} \psi_i \psi_{i+r}^*$$

power sums / KP 'time parameters'

Boson-Fermion correspondence

SUMMARY

- The character table of symmetric groups yields the change of basis $p_\mu \rightarrow s_\lambda$, $\lambda, \mu \vdash n$

- s_λ is a basis distinguished by positivity

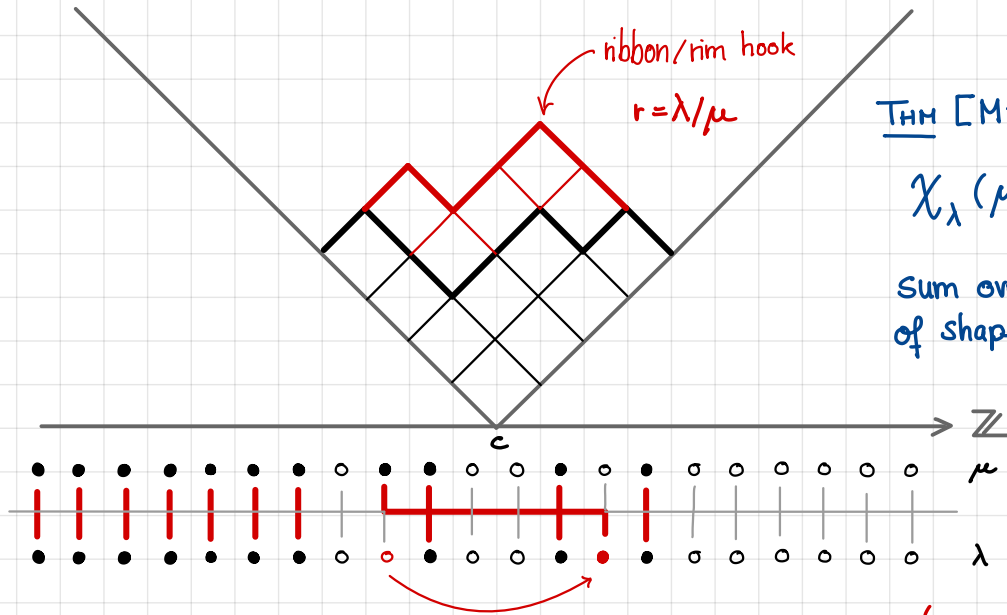
$c_{\lambda\mu}^\nu \in \mathbb{Z}_{\geq 0} \rightsquigarrow$ intersection cohomology $H^*(Gr_{k,n})$

Q. How does one compute $\chi_\lambda(\mu)$?

A. Murnaghan-Nakayama rule

(combinatorial algorithm using ribbons / rim hooks)

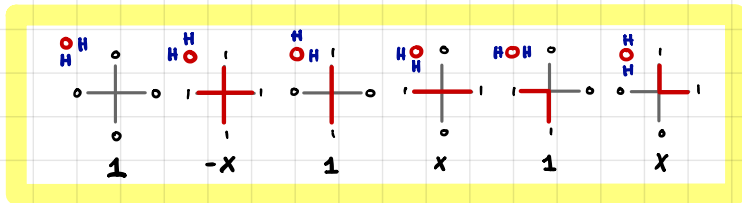
THE MURNAGHAN-NAKAYAMA RULE



TM [M-N]

$$\chi_\lambda(\mu) = \sum_{T \in \mathcal{T}} \prod (-1)^{ht(r)-1}$$

sum over all rim hook tableaux of shape λ and weight μ .



'special' configurations of the asymmetric 6-vertex model: only one pair of allowed!

§2. The quantum characteristic map

Reminder: the Iwahori-Hecke algebra of type A

DEF $H_n = H_n(t)$ is the $\mathbb{C}(t)$ -algebra with generators $T_i, i=1, \dots, n-1$ and relations

$$T_i^2 + (1-t)T_i - t = 0, \quad T_i T_{i+1} T_i = T_{i+1} T_i T_{i+1}, \quad T_i T_j = T_j T_i \quad (|i-j| > 1)$$

Grothendieck
ring of Hecke
algebras

$$\mathcal{R}_t = \bigoplus_{n \in \mathbb{N}} \mathcal{R}_t^n \xrightarrow{\sim} \mathcal{Z} = \bigoplus_{n \in \mathbb{N}} \mathcal{Z}(H_n) \quad \text{centre of } H_n(t)$$

ch_t
quantum
characteristic
map

F_t

quantum
Frobenius [Wan-Wang]
map

$$\text{Sym}_t = \bigoplus_{n \in \mathbb{N}} \text{Sym}_t^n = \text{Sym}^n \otimes_{\mathbb{C}} \mathbb{C}(t)$$

ring of symmetric fctns

Plethystic variable substitution

THM [Wan-Wang] (inverse) 'quantum Frobenius formula'

$$s_\lambda[X] = \sum_{\mu} \chi_t^\lambda(T_{w_\mu}) (t-1)^{\ell(\mu)} m_\mu\left[\frac{X}{t-1}\right]$$

Schur function in alphabet X
 irreducible Hecke character $\lambda \vdash n$
 $T_{w_\mu} = T_{i_1} \dots T_{i_\ell}$ where $w_\mu = s_{i_1} \dots s_{i_\ell}$ reduced and of cycle type μ
 monomial symmetric function in alphabet $\frac{X}{t-1}$

Equivalently,

$$p_r\left[\frac{X}{t-1}\right] = p_r[X] / (t-1)^r$$

$$s_\lambda[(t-1)X] = \sum_{\mu} \chi_t^\lambda(T_{w_\mu}) (t-1)^{\ell(\mu)} m_\mu[X]$$

'dual' Schur function w.r.t. $\langle P_\lambda, Q_\mu \rangle_t = \delta_{\lambda\mu}$
HL-functions

A Hecke version of the BF correspondence

Fermionic Fock space: $\Lambda^{\frac{\infty}{2}, c} V_t = \Lambda^{\frac{\infty}{2}, c} V \otimes_{\mathbb{C}} \mathbb{C}(t)$

'free' six-vertex transfer matrix

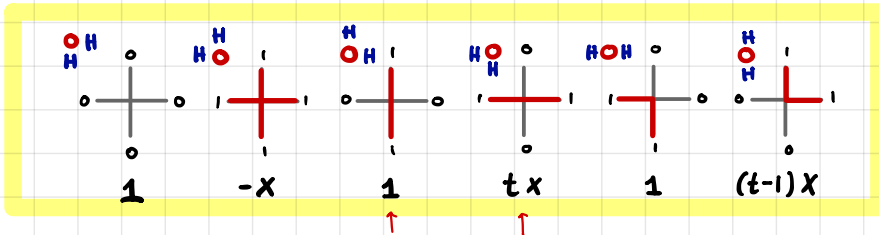
PROP (CK) $\sigma_{\lambda, c} \mapsto \langle \sigma_{\lambda, c}, \prod_{i \geq 1} A(x_i; t) \sigma_{\emptyset, c} \rangle = S_{\lambda}[(t-1)X]$

Maya diagram $\hat{=}$ spin configuration

$\langle \sigma', A(x) \sigma \rangle = \sum_{\epsilon_i} \dots$

$A(x; t) = e^{H_t[p]}$

$H_t[p] = \sum_{r > 0} \frac{(t^r - 1)}{r} p_r \sum_{i \in \mathbb{Z}} \psi_i \psi_{i+r}^*$

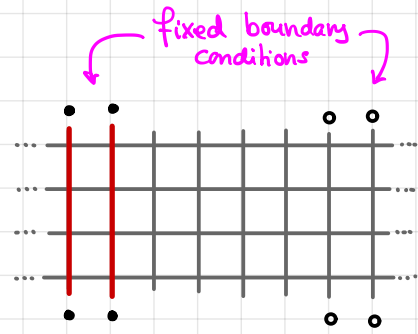
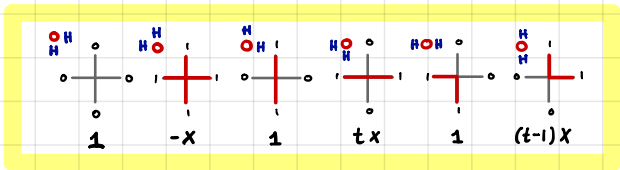
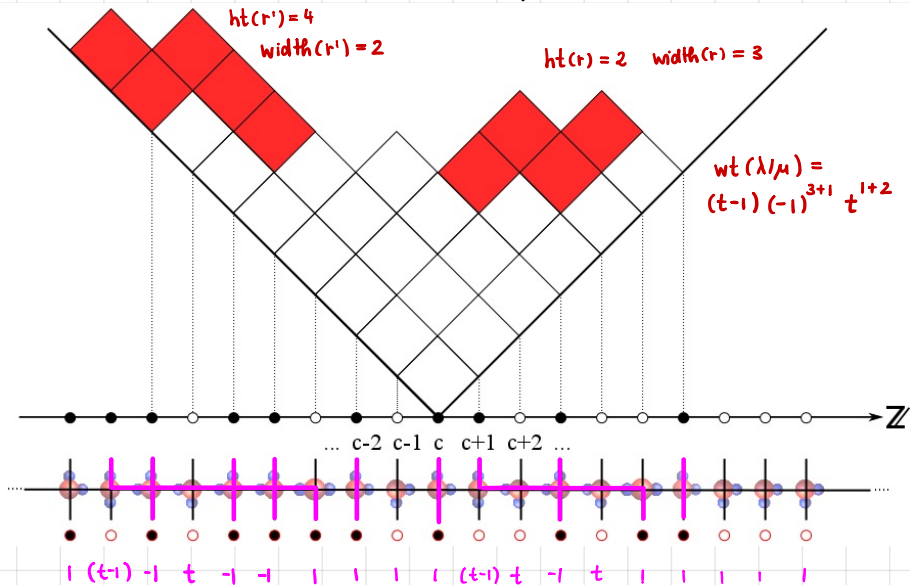


six vertex configurations & Boltzmann weights

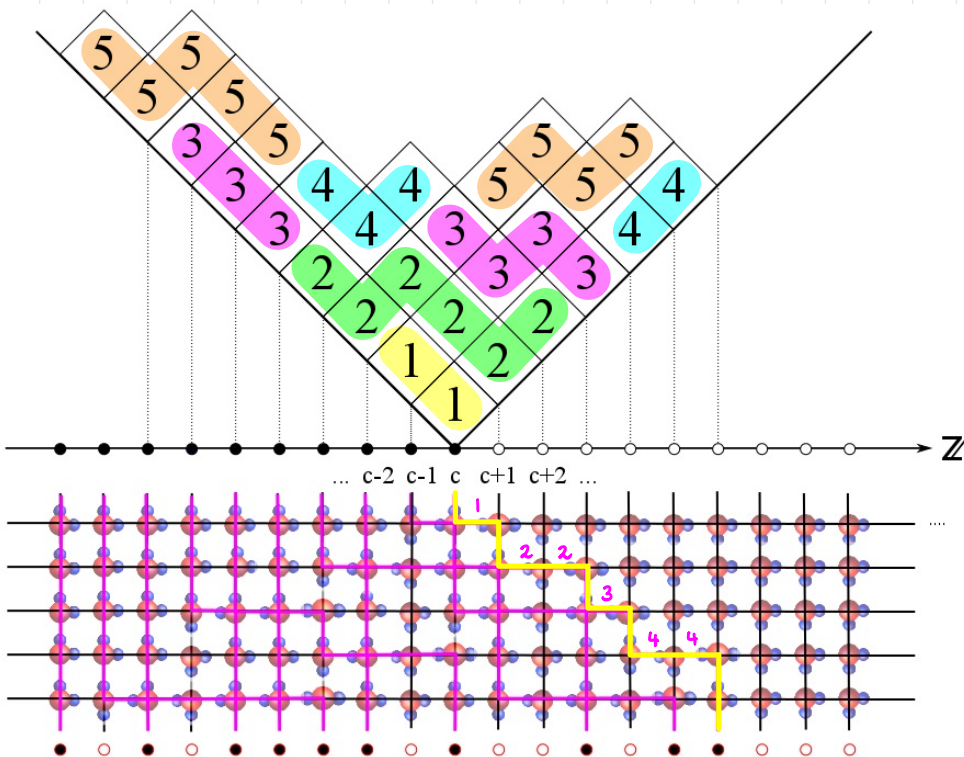
Rmk These are different from the weights for Tokuyama ice [Brubaker-Bump-Friedberg] and different boundary conditions

Broken rim hooks & asymmetric ice configurations

$$wt(\lambda/\mu) = (t-1)^{\#\text{rim hooks}_r} \prod_{r \geq 1} (-1)^{ht(r)-1} t^{\text{width}(r)-1}$$



Broken rim hook tableaux and Hecke characters



THM [Ram]

Hecke characters are given as weighted sums over broken rim hook tableaux.

$$\chi_t^\lambda (T_{w_\mu}) = \sum_{\mathcal{T}} \text{wt}(\mathcal{T})$$

$$\text{wt}(\mathcal{T}) = \prod_i \text{wt}(\lambda^{(i)} / \lambda^{(i+1)})$$

w_μ representative of conjugacy class fixed by μ

PROP The 6-vertex lattice configurations are in bijection with broken rim hook tableaux.

§3. Vertex operators & transfer matrices

DEF (JING) Define the following (bosonic) vertex operators $\Phi^\pm(x;t): \text{Sym}_t \longrightarrow \text{Sym}_t$

$$\Phi^-(x;t) = e^{\sum_{r>0} \frac{1-t^r}{r} p_r x^r} e^{\sum_{r>0} (t^r-1) \frac{\partial}{\partial p_r} x^{-r}}, \quad \Phi^+(x;t) = e^{\sum_{r>0} \frac{t^r-1}{r} p_r x^r} e^{\sum_{r>0} (1-t^r) \frac{\partial}{\partial p_r} x^{-r}}$$

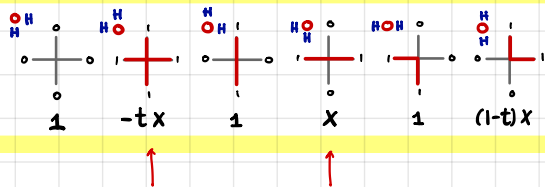
THM (JING) $\Phi^-(x;t) = \sum_{r \in \mathbb{Z}} \Phi_r^- x^r$, $\Phi_{\lambda_1}^- \dots \Phi_{\lambda_\ell}^- 1 = Q_\lambda$ Hall-Littlewood function

Six-vertex construction

THM (CK) Let $z_t: \Lambda^{\frac{\infty}{2},c} V_t \rightarrow \text{Sym}_t$ be the t -extension of the BF-correspondence, then

$$z_t \circ A^{-1}(x) A^*(x) = \Phi^-(x;t) \circ z_t \quad \text{and} \quad z_t \circ A(x) (A^{-1}(x))^* = \Phi^+(x;t) \circ z_t$$

where $A^{-1}(x;t) = A(tx;t^{-1})$ is the six-vertex transfer matrix with weights



'fermionic' construction
of Hall-Littlewood
functions

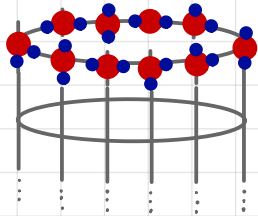
q-boson construction:
[Tsilchenik '06]
C.K. CMP318(2013)173

LLT polynomials [Brubaker, Buciumas, Bump, Gustafsson]

§4. Periodic boundary conditions

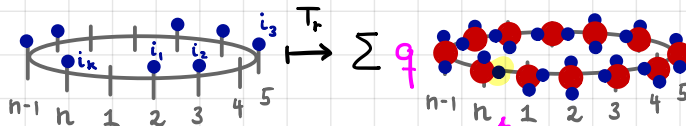
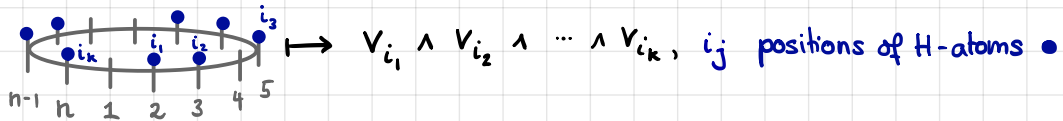
$$\bigwedge_{\mathbb{Z}^d, k} V \otimes \mathbb{C}(t) \rightarrow \bigwedge \mathbb{C}^n(t) = \bigoplus_{k=0}^n \bigwedge^k \mathbb{C}^n(t)$$

∞ wedge finite wedge



Denote by $T(x)$ the six-vertex transfer matrix with quasi-periodic boundary conditions and twist parameter q

$$T(x) = \sum_{r=0}^n x^r T_r, \quad T_r : \bigoplus_{d \in \mathbb{Z}} q^d \otimes \bigwedge \mathbb{C}^n(t) \rightarrow \bigoplus_{d \in \mathbb{Z}} q^d \otimes \bigwedge \mathbb{C}^n(t)$$

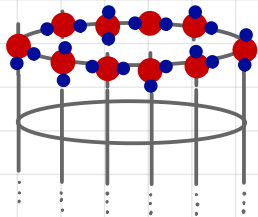


H-atom between site n and site 1 results in factor q !

§4. Periodic boundary conditions

$$\Lambda^{\infty, k} V \otimes \mathbb{C}(t) \rightarrow \Lambda \mathbb{C}^n(t) = \bigoplus_{k=0}^n \Lambda^k \mathbb{C}^n(t)$$

∞ wedge finite wedge



Denote by $T(x)$ the six-vertex transfer matrix with quasi-periodic boundary conditions and twist parameter q

$$T(x) = \sum_{r=0}^n x^r T_r, \quad T_r : \bigoplus_{d \in \mathbb{Z}} q^d \otimes \Lambda^k \mathbb{C}^n(t) \rightarrow \bigoplus_{d \in \mathbb{Z}} q^d \otimes \Lambda^k \mathbb{C}^n(t)$$



$\mapsto V_{i_1} \wedge V_{i_2} \wedge \dots \wedge V_{i_k}, \quad i_j \text{ positions of H-atoms}$

Bethe ansatz: $T(x)$ is diagonalizable with eigenvalues

$$\tau(x) = \underbrace{(1 + (-1)^k q x^n t^n)}_{\text{Bethe roots}} \prod_{i=1}^k \frac{1 - y_i x}{1 - t y_i x}, \quad y_i^n = (-1)^{k-1} q$$

Define a 'renormalised' transfer matrix $H(x) : \bigoplus_{d \in \mathbb{Z}} q^d \otimes \Lambda^k \mathbb{C}^n(t) \rightarrow \bigoplus_{d \in \mathbb{Z}} q^d \otimes \Lambda^k \mathbb{C}^n(t)$

$$H(x) = \frac{T(x)}{\underbrace{1 + (-1)^k q x^n t^n}} = \sum_{r \geq 0} x^r H_r$$

Cylindric Hecke characters

renormalised 6-vertex transfer matrix

DEF $q^d \chi_t^{\lambda/d/\mu}(v) (t-1)^e = \langle v_\lambda, H_{v_1} \dots H_{v_e} v_\mu \rangle,$

$$H(x) = \sum_{r \geq 0} x^r H_r$$

$$d = \frac{|\mu| + |v| - |\lambda|}{n}$$

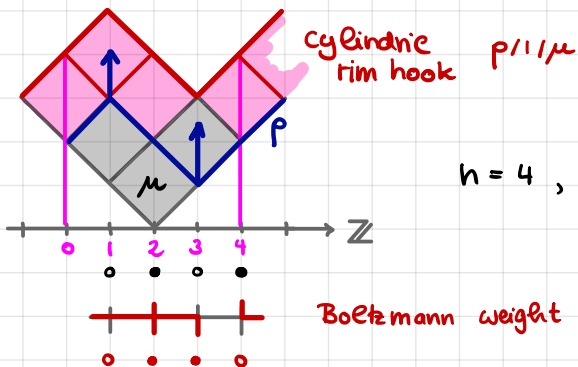
LEMMA (Cylindric Murnaghan-Nakayama rule)

(i) $\chi_t^{\lambda/d/\mu}(\dots, v_i + n, \dots) = (-1)^{k-1} t^n \chi_t^{\lambda/d-11\mu}(\dots, v_i, \dots)$

(ii) $\chi_t^{\lambda/d/\mu}(v, r) = \sum_p \chi_t^{\lambda/d/p}(v) \chi_t^{p/\mu}(r) + \sum_p \chi_t^{\lambda/d-11p}(v) \chi_t^{p/11\mu}(r)$

The rules (i), (ii) allow one to compute $\chi_t^{\lambda/d/\mu}$.

Cylindric part



$$h = 4, k = 2 \quad \chi_t^{(2,0)/1/(2,1)}(3) = (-1)^{2-1} t^{2-1}$$

The coalgebra of cylindric Hecke characters

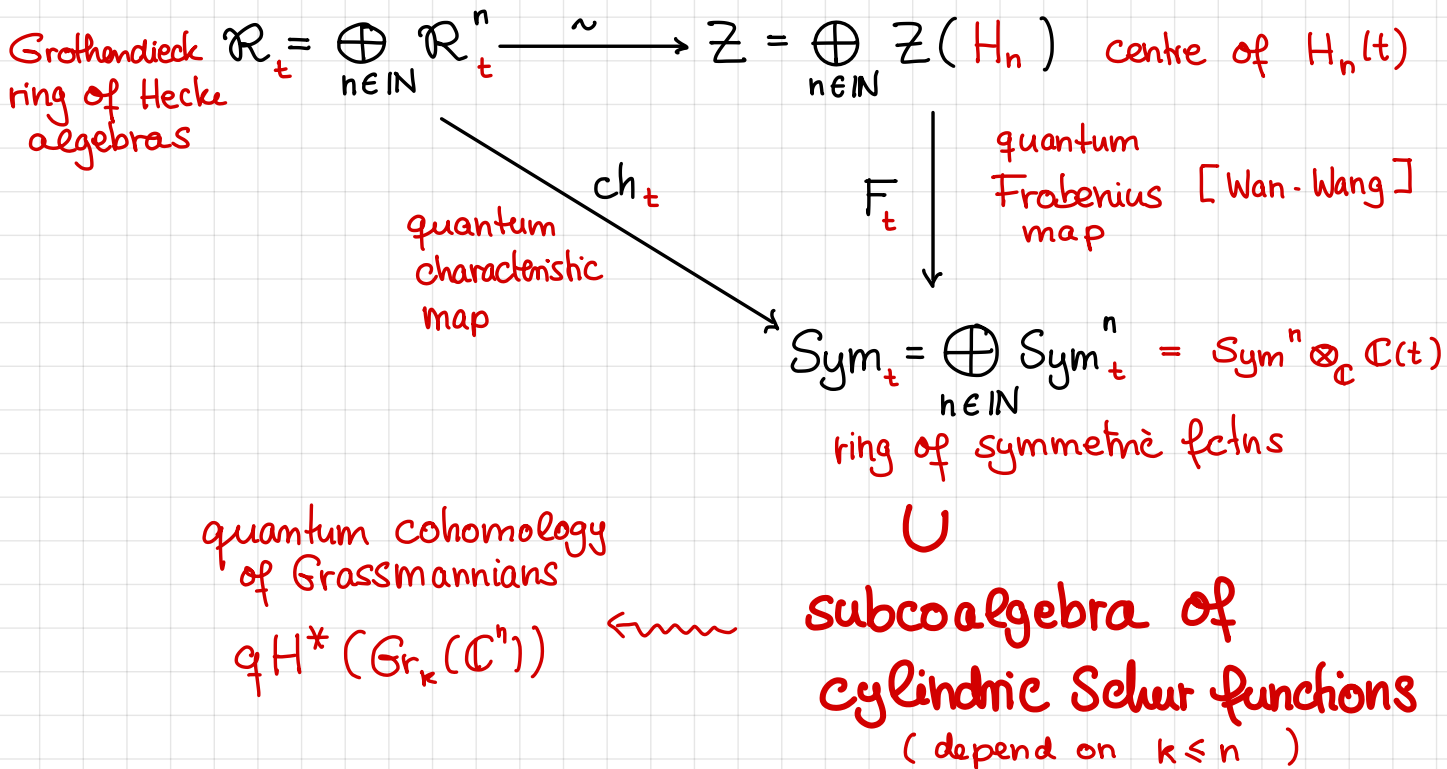
Main Theorem (C.K.) The cylindric Hecke characters $\chi_t^{\lambda/d|\emptyset}$ with $\lambda \in \square_{n-k}^k$, $d \geq 0$ span an ∞ -dimn'l subcoalgebra of \mathcal{R}_t (Grothendieck ring) with

$$\text{Res}_{H_{m'} \otimes H_{m''}}^{H_{m'+m''}} \chi_t^{\lambda/d|\emptyset} = \sum_{d'+d'' \leq d} \sum_{\substack{\mu \vdash m' \\ \nu \vdash m''}} C_{\mu\nu}^{\lambda, d-d'-d''} \chi_t^{\mu/d'|\emptyset} \otimes \chi_t^{\nu/d''|\emptyset}$$

where $C_{\mu\nu}^{\lambda, d} \in \mathbb{Z}_{\geq 0}$ are the 3-point genus 0 Gromov-Witten invariants of the Grassmannian $\text{Gr}_k(\mathbb{C}^n)$.

For $d=0$ one recovers $H^*(\text{Gr}_k(\mathbb{C}^n))$ as coalgebra with $C_{\mu\nu}^{\lambda, 0} = c_{\mu\nu}^{\lambda}$ being the Littlewood-Richardson coefficients. (intersection cohomology)

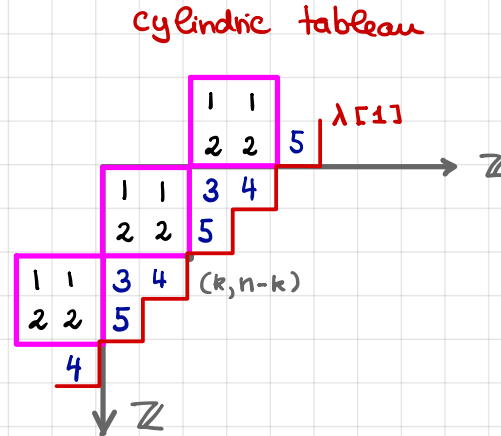
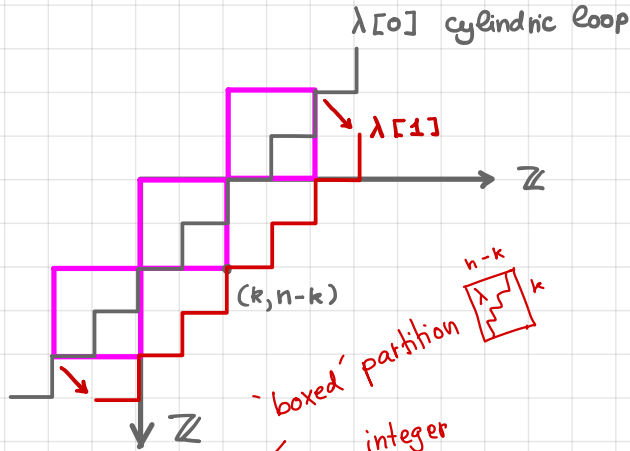
Q. What is the image of cylindric Hecke characters under Macdonald's characteristic map?



Combinatorial definition of cylindric Schur functions

[Gessel-Krotenthaler, Postnikov, Lam, McNamara ...] [D. Palazzo, C.K.]

↳ positive coalgebra



DEF $s_{\lambda/d/\phi} = \sum_T x^T$, T cylindric tableau of shape $\lambda[d]$

THM (C.K., D. Palazzo) Cylindric Schur functions form a positive subcoalgebra $\subset \text{Sym}$

A cylindric (ae) Boson-Fermion correspondence

$$\mathcal{R}_t = \bigoplus_{m \geq 0} \mathcal{R}_t^m \xrightarrow{\sim} \mathcal{Z}_t = \bigoplus_{m \geq 0} \mathcal{Z}(\mathcal{H}_m(t)) \text{ centres of Hecke algebras}$$

$\downarrow s$ $\downarrow F_t$ Frobenius map

BF correspondence

$$\Lambda_{\mathbb{Z}, \mathbb{C}}^{\infty, \mathbb{C}} V \otimes \mathbb{C}(t) \longrightarrow \Lambda_t = \Lambda \otimes_{\mathbb{C}} \mathbb{C}(t)$$

$\downarrow c = k \bmod n$ \downarrow rim hook algorithm

Satake correspondence

$$\Lambda^k \mathbb{C}^n(t) \otimes \mathbb{C}[q] \longrightarrow qH^*(G_r(\mathbb{C}^n)) \otimes \mathbb{C}(t), \quad qH^*(G_r(\mathbb{C}^n)) = \mathbb{C}[q, e_1, \dots, e_k, h_1, \dots, h_{n-k}] / \sim$$

Bethe ansatz \Leftrightarrow $\left(\sum_{i=0}^k (-x)^i e_i \right) \left(\sum_{j=0}^{n-k} x^j h_j \right) = 1 + q(-1)^k x^n$
 [C.K., Stroppel]

cylindric boson-fermion correspondence

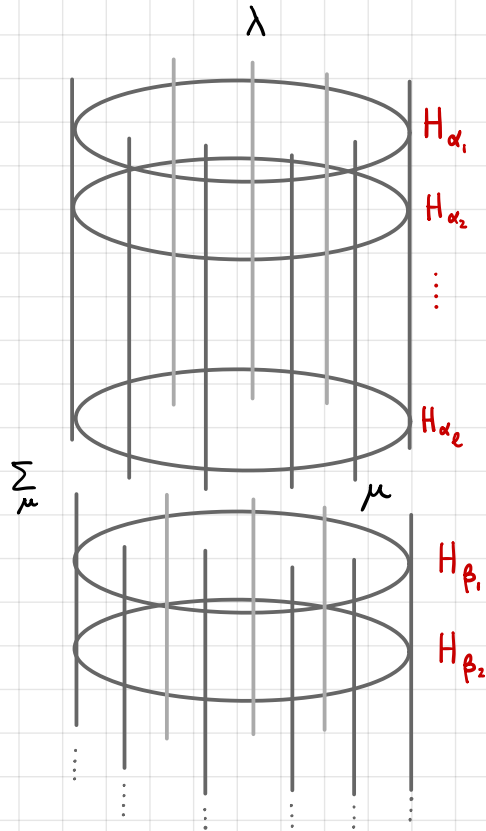
$$v_\lambda = v_{\lambda_1 + k-1} \wedge \dots \wedge v_{\lambda_k} \mapsto \langle v_\lambda, H(x_1)H(x_2) \dots v_\emptyset \rangle = \sum_{d \geq 0} q^d s_{\lambda'/d/\emptyset} [(t-1)X] \text{ cylindric Schur functions}$$

$$= \sum_{d \geq 0} q^d \sum_{\alpha} \chi_t^{\lambda'/d/\emptyset}(\alpha) (t-1)^{e(\alpha)} m_\alpha [X]$$

cylindric Hecke characters

Proof (Sketch)

Use bra-ket notation: $\langle \lambda | \otimes | \mu \rangle \equiv \langle \nu_\lambda, \otimes \nu_\mu \rangle$



$$q^d \chi_t^{\lambda/d|\phi}(\alpha, \beta) (t^{-1})^{e(\alpha)+e(\beta)} = \langle \lambda | H_\alpha H_\beta | \phi \rangle$$

$$= \sum_{\mu} \langle \lambda | H_\alpha | \mu \rangle \langle \mu | H_\beta | \phi \rangle$$

$$= q^d (t^{-1})^{e(\alpha)+e(\beta)} \sum_{d'=0}^d \sum_{\mu} \underbrace{\chi_t^{\lambda/d-d'|\mu}(\alpha)}_{\text{skew cylindric Hecke character}} \chi_t^{\mu/d'|\phi}(\beta)$$

skew cylindric Hecke character

Use the Bethe ansatz (eigenbasis of H_α)

and Bertam-Vafa-Intrigatator formula for $C_{\mu\nu}^{\lambda,d}$

to show that
$$\chi_t^{\lambda/d|\mu} = \sum_{d'=0}^d \sum_{\nu} C_{\mu\nu}^{\lambda,d-d'} \chi_t^{\nu/d'|\phi}$$

Proof (Sketch)

$$q^d \chi_t^{\lambda/d|\phi}(\alpha, \beta) (t-1)^{\ell(\alpha)+\ell(\beta)} = \langle \lambda | H_\alpha H_\beta | \phi \rangle = \sum_{\mu} \langle \lambda | H_\alpha | \mu \rangle \langle \mu | H_\beta | \phi \rangle$$

$$= q^d (t-1)^{\ell(\alpha)+\ell(\beta)} \sum_{d'=0}^d \sum_{\mu} \chi_t^{\lambda/d-d'|\mu}(\alpha) \chi_t^{\mu/d'|\phi}(\beta)$$

Employing eigenbasis of the H_r (Bethe ansatz)

$$q^d \chi_t^{\lambda/d|\mu}(\alpha) (t-1)^{\ell(\alpha)} = \sum_{\xi} \langle \lambda | H_\alpha | \xi \rangle \langle \xi | \mu \rangle, \quad \xi_i^n = (-1)^{k-1} q$$

Bethe vectors Bethe roots

$$= \sum_{\xi} h_\alpha [(t-1)\xi] \frac{s_\lambda(\xi) s_\mu(\xi^{-1})}{n^k \prod_{i < j} |\xi_i - \xi_j|^{-2}}$$

Bethe formula for cylindric Hecke characters

$$\sum_{\nu} \frac{s_\nu(\xi^{-1}) s_\nu(\xi)}{n^k \prod_{i < j} |\xi_i - \xi_j|^{-2}} = \delta_{\xi \tilde{\xi}}$$

$$= \sum_{\nu} \left(\underbrace{\sum_{\xi} \frac{s_\lambda(\xi) s_\mu(\xi^{-1})}{n^k \prod_{i < j} |\xi_i - \xi_j|^{-2}}}_{q^{d-d'} C_{\mu\nu}^{\lambda, d-d'}} s_\nu(\xi^{-1}) \right) \underbrace{\sum_{\xi} h_\alpha [(t-1)\xi] \frac{s_\nu(\xi)}{n^k \prod_{i < j} |\xi_i - \xi_j|^{-2}}}_{q^{d'} \chi_t^{\nu/d'|\phi}(\alpha)}$$

[Bertam-Vafa-Intelligator]

THANK YOU FOR
YOUR ATTENTION!